Effects of network topology, transmission delays, and refractoriness on the response of coupled excitable systems to a stochastic stimulus  Daniel B. Larremore,					

In a recent report,<sup>2</sup> we presented an analysis of the

## **B.** Model dynamics

By considering a large ensemble of realizations of the above stochastic process on the same network, we can define the probability that node i is at state  $x_i^t$  at time t as  $p_i^t \delta x P$ . The probabilities  $p_i^t$  evolve in one time step by

(4)

$$p_i^{t \triangleright 1} \delta m_i \triangleright \frac{1}{4} p_i^t \delta m_i \qquad 1 \triangleright, \tag{5}$$

and we also have the normalization condition

$$p_i^{\dagger} \tilde{0} 0 \tilde{P} \stackrel{\mathcal{M}_i}{\cancel{4}} 1 \qquad p_i^{\dagger} \tilde{0} j \tilde{P}, \tag{6}$$

where  $r_i^t$  in Eq. (2) is the rate of transitions from the ready to the excited state, given by

where  $l_j^t$  is one if node j is excited at time t and zero otherwise, and E½ denotes an ensemble average. Assuming that the neighbors of node i being excited are independent events, we obtain, letting  $p_i^t \tilde{0} 1 P$   $p_i^t$ ,

Now, we use the fact that the largest eigenvalue of A,  $\lambda$ , is

determines  $\bar{p}_i$ ) and decays exponentially with the number of expected excitations from its neighbors. In terms of the aggregate response  $\hat{F}$ , Eq. (22) becomes, after multiplying by  $A_{ki}$ , summing over k and i, and normalizing,

$$\frac{\mathrm{d}\hat{F}}{\mathrm{d}\eta} \, \frac{\mathrm{hd}^{\mathrm{out}} \bar{p}^2 \mathrm{e}^{-\mathrm{A}\bar{p}} \mathrm{i}}{\mathrm{hd}\, \mathrm{i}} \, . \tag{23}$$

## C. Dynamics near the critical regime

More precisely, the value of stimulus  $\eta_{low}$  ( $\eta_{high}$ ) corresponding to a low (high) threshold of activity  $\hat{F}_{low}$  ( $\hat{F}_{high}$ ) are found and the dynamic range is calculated as

Using our approximations to the response  $\hat{\mathsf{F}}$  as a function of stimulus  $\eta$ 

and then connecting them using the configuration model. <sup>29</sup> In some cases, an additional fourth step was used to change the assortativity coefficient  $\rho$ , defined in Eq. (30), of a critical (i.e., with  $\lambda$  ¼ 1) scale-free network, making this network more assortative (disassortative) by choosing two links at random, and swapping their destination connections only if the resulting swap would increase (decrease)  $\rho$ . This swapping allows for the degree of assortativity (and thereby,  $\lambda$ ) to be modified while preserving the network's degree distribution. <sup>8,19</sup>

## B. Results of numerical experiments

We first demonstrate the ability of the non-perturbative approximation to predict aggregate network behavior in a variety of conditions. Fig. 2 shows a multitude of simulations (symbols) with the predicted behavior of Eq. (17) overlaid (lines). The cases considered in Fig. 2 include different combinations of topology, assortativity, largest eigenvalue  $\lambda$ , delays, and number of refractory states. The number of refractory states  $m_i$  was chosen either constant,  $m_i$  ¼ m, or randomly chosen with equal probability among  $f1, 2, ..., m_{max}g$ 

of the refractory states. Equation (21) predicts that the response should scale as hm  $\triangleright 1/2i^{-1}$ . The inset shows how, after multiplication by hm  $\triangleright 1/2i$ , the response curves collapse into a single curve. Figure 4 also depicts a linear relationship,  $\hat{F} = \delta \lambda - 1 \triangleright$  for  $\lambda > 1$ . Making a connection with the theory of nonequilibrium phase transitions in which  $\hat{F} = \delta \lambda - \lambda_c \flat^\beta$ , we derive  $\lambda_c$  ¼ 1 and the critical exponent  $\beta$  ¼ 1.

Figure 5 shows the response  $\hat{F}$  close to  $\eta$  ¼ 1 calculated for various values of m from the simulation (symbols), and from Eq. (23) (solid lines). Equation (23) describes well the slope of  $\hat{F}$  close to  $\eta$  ¼ 1. An important observation is that as m grows, the relative slope  $\hat{F}$   $^{1}d\hat{F}/d\eta$  at  $\eta$  ¼ 1 decreases. Therefore, if the typical refractory period m is large, the response  $\hat{F}$  saturates [e.g., reaching 90% of  $\hat{F}$ 01 $^{1}$ 1 for smaller values of  $\eta$ .

Transmission delays, as in the analogous system of gene regulatory networks, <sup>18</sup>

determined from the distribution of delays. In Fig. 6, we					

